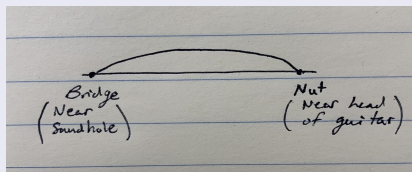


Derivation of the Wave Equation

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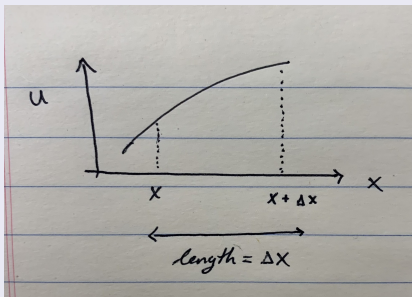
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- Consider a string subject to the conditions:
 - Only vibrates in a fixed plane
 - Constant mass density ρ
 - Perfectly elastic and experiences no resistance
 - Transverse vibrations are small and occur in the x-axis (overall motion of the string depends only on the y-displacement and time)
- Such assumptions are not unrealistic. Consider a taut guitar string:



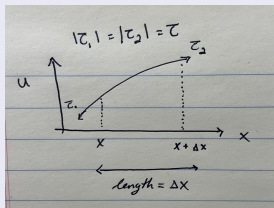
Closeup of String Segment

- Consider an infinitesimal region on the string of length Δx :

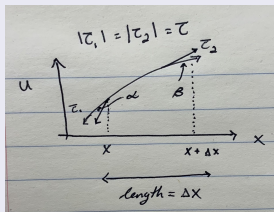


Relevant Forces

- String is taut \implies tension force τ :

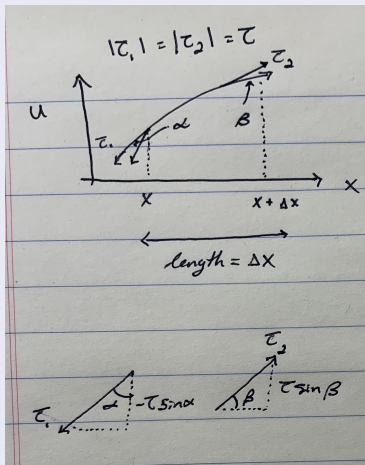


- Magnitude of τ must be equal on both ends but the positions we consider differ in angle \implies two tension components τ_1 and τ_2 at angles α and β :



Relevant Forces Continued

- No motion in x-direction so we only consider the y-components of these tension forces:



Considering the Trigonometry of the String Segment

- Tension forces must balance because $|\tau_1| = |\tau_2| = \tau$ so:

$$\tau \sin(\beta) = \tau \sin(\alpha)$$

- By Newton's Second Law, the sum of tension forces must be equal to a mass multiplied by an acceleration:

$$-\tau \sin(\alpha) + \tau \sin(\beta) = F = ma$$

- Consider the small angle approximations $\sin(\theta) \approx \theta$, $\tan(\theta) \approx \theta$.
 $\sin(\theta) = \tan(\theta)$ under this approximation so make the replacement:

$$-\tau \tan(\alpha) + \tau \tan(\beta) = F = ma$$

Applying Newton's Second Law

- Recall in 1 dimension that:

$$F = ma = m \frac{d^2x}{dt^2}$$

- Because string depends on y-position and time $u(y, t)$:

$$F = ma = m \frac{\partial^2 u}{\partial t^2}$$

- Combining this with the previous trig replacement:

$$-\tau \tan(\alpha) + \tau \tan(\beta) = F = m \frac{\partial^2 u}{\partial t^2}$$

Calc 1 Definition of Tangents

- Recall from Calc 1 that a tangent line to a point is considered its derivative and the limit definition of a derivative is:

$$\lim_{x \rightarrow 0} \frac{f(x + \Delta x) - f(x)}{\Delta x} = \frac{df}{dx}$$

- In this multivariable case $u(y, t)$, the variable in question here is the y -component so the tangent becomes a partial derivative
- Also remember that we consider a segment of string so evaluate these partials at both ends:

$$\tau \tan(\alpha) = \tau \left. \frac{\partial u}{\partial x} \right|_{(x)}$$

$$\tau \tan(\beta) = \tau \left. \frac{\partial u}{\partial x} \right|_{(x+\Delta x)}$$

Applying to Earlier Expression

- From original expression:

$$-\tau \tan(\alpha) + \tau \tan(\beta) = F = m \frac{\partial^2 u}{\partial t^2}$$

- Replace definitions for $\tan(\alpha)$ and $\tan(\beta)$ from last slide:

$$-\tau \left. \frac{\partial u}{\partial x} \right|_{(x+\Delta x)} + \tau \left. \frac{\partial u}{\partial x} \right|_{(x+\Delta x)} = F = m \frac{\partial^2 u}{\partial t^2}$$

- Rearranging to get to a familiar form:

$$\tau \left(\left. \frac{\partial u}{\partial x} \right|_{(x+\Delta x)} - \left. \frac{\partial u}{\partial x} \right|_{(x+\Delta x)} \right) = m \frac{\partial^2 u}{\partial t^2}$$

Rearranging

- Notice that left side of the expression:

$$\tau \left(\frac{\partial u}{\partial x} \Big|_{(x+\Delta x)} - \frac{\partial u}{\partial x} \Big|_{(x)} \right) = m \frac{\partial^2 u}{\partial t^2}$$

is again in the form for the definition of a derivative:

$$\lim_{\Delta x \rightarrow 0} \frac{f(x + \Delta x) - f(x)}{\Delta x} = \frac{df}{dx}$$

without the Δx term in the denominator. Divide both sides by this Δx term so left side is turned into a second derivative:

$$\tau \frac{\partial^2 u}{\partial x^2} = \frac{m}{\Delta x} \frac{\partial^2 u}{\partial t^2}$$

- Again, we have a multivariable expression so the derivative is a partial here as well.

Wave Equation Recovered

- Recall we are considering an infinitesimal of string and currently have:

$$\tau \frac{\partial^2 u}{\partial x^2} = \frac{m}{\Delta x} \frac{\partial^2 u}{\partial t^2}$$

- Where m refers to the overall mass of the string and Δx refers to the length of this small segment.
- Define $\frac{m}{\Delta x}$ to be the mass density ρ of the string.
- Divide τ to right side. $\frac{\tau}{\rho}$ has units of $\frac{[\text{length}]^2}{[\text{time}]^2}$ so call this c^2 (a velocity).

$$\frac{\partial^2 u}{\partial x^2} = \frac{1}{c^2} \frac{\partial^2 u}{\partial t^2}$$

- Leaving you with the standard form of the wave equation [1] in one spatial dimension.
- Note that this could be easily extended to multiple dimensions since the system is linear.

- [1] Nakhlé H Asmar. *Partial Differential Equations and Boundary Value Problems*. Upper Saddle River, N.J.: Prentice Hall, 2000.